

## Clarification of Expectation Values

Expectation value is the ensemble average over many identical systems. Does not have to correspond to an actual measured value.

Ex: Spin- $1/2$  particles can have spin-up  $\frac{\hbar}{2}|\uparrow\rangle$  or spin-down  $-\frac{\hbar}{2}|\downarrow\rangle$ .

Expectation value  $\Rightarrow 0$ , but not a possible measurement

\* This was assuming  $\hat{S}_z |\uparrow\rangle_z = +\frac{\hbar}{2} |\uparrow\rangle_z$

$$\hat{S}_z |\downarrow\rangle_z = -\frac{\hbar}{2} |\downarrow\rangle_z$$

We are going to prove that in the basis of an operator  $\hat{Q}$ , a state  $|\Psi\rangle$  has expectation value  $\langle\Psi|\hat{Q}|\Psi\rangle = \sum_n q_n |c_n|^2$ .

$q_n$  are eigenvalues of  $\hat{Q}$ ,  $c_n$  are coefficients.

Q1: Show  $|\Psi\rangle = \sum_n c_n |q_n\rangle$  using bra-ket notation.

A1:  $|\Psi\rangle = \sum_n |q_n\rangle \langle q_n | \Psi \rangle = \sum_n c_n |q_n\rangle$  where  $c_n \equiv \langle q_n | \Psi \rangle$   
 $\hat{Q} |q_n\rangle = q_n |q_n\rangle$

Q2: Use Q1 to show  $\langle\Psi|\hat{Q}|\Psi\rangle = \sum_n q_n |c_n|^2$

A2:  $\langle\Psi|\hat{Q}|\Psi\rangle = \langle\Psi|\hat{Q} \sum_n c_n |q_n\rangle = \sum_n \langle\Psi|\hat{Q}|q_n\rangle$   
 $= \sum_n c_n \langle\Psi|q_n|q_n\rangle = \sum_n q_n c_n \langle\Psi|q_n\rangle$   
 $= \sum_n q_n c_n c_n^* = \sum_n q_n |c_n|^2$  since  $c_n^* = \langle\Psi|q_n\rangle$

Q3: Use the identity  $\mathbb{1}$  expressed in terms of the eigenstates of  $\hat{Q}$  to show  $\langle\Psi|\hat{Q}|\Psi\rangle = \sum_n q_n |c_n|^2$ .

A3:  $\langle\Psi|\hat{Q}|\Psi\rangle = \langle\Psi|\hat{Q}\mathbb{1}|\Psi\rangle = \langle\Psi|\hat{Q} \sum_n |q_n\rangle \langle q_n | \Psi \rangle$   
 $= \sum_n \langle\Psi|q_n|q_n\rangle c_n = \sum_n q_n c_n^* c_n = \sum_n q_n |c_n|^2$

Q4: What if  $\hat{Q}$  has a continuous spectrum?  $\hat{Q}|q\rangle = q|q\rangle$ ?  
 Use the identity  $\mathbb{I}$  in continuous form w.r.t.  $q$  to find the expectation value  $\langle \Psi | \hat{Q} | \Psi \rangle$ .

A4:  $\mathbb{I} = \int_{-\infty}^{\infty} |q\rangle \langle q| dq$

$$\begin{aligned} \langle \Psi | \hat{Q} | \Psi \rangle &= \langle \Psi | \hat{Q} \mathbb{I} | \Psi \rangle = \langle \Psi | \hat{Q} \int_{-\infty}^{\infty} |q\rangle \langle q| dq | \Psi \rangle \\ &= \int_{-\infty}^{\infty} \langle \Psi | \hat{Q} | q \rangle \langle q | \Psi \rangle dq \\ &= \int_{-\infty}^{\infty} \langle \Psi | q | q \rangle \langle q | \Psi \rangle dq = \int_{-\infty}^{\infty} q \underbrace{\langle \Psi | q \rangle \langle q | \Psi \rangle}_{|\Psi|^2} dq \\ \Rightarrow \langle \Psi | \hat{Q} | \Psi \rangle &= \int_{-\infty}^{\infty} q |\Psi|^2 dq \end{aligned}$$

Q5: What if  $|\Psi\rangle$  isn't expanded in the basis of  $\hat{Q}$ ?  
 Show  $\langle \Psi | \hat{Q} | \Psi \rangle = \sum_{n,m,i,j} c_m^* c_n Q_{ij} \delta_{mj} \delta_{in}$ .

A5:  $\langle \Psi | \hat{Q} | \Psi \rangle = \langle \Psi | \hat{Q} \sum_n |q_n\rangle \langle q_n | \Psi \rangle$        $c_n \equiv \langle q_n | \Psi \rangle$   
 $= \sum_n \langle \Psi | \hat{Q} | q_n \rangle \langle q_n | \Psi \rangle$        $c_m^* \equiv \langle \Psi | q_m \rangle$   
 $= \sum_{n,m} \langle \Psi | q_m \rangle \langle q_m | \hat{Q} | q_n \rangle \langle q_n | \Psi \rangle$   
 $= \sum_{n,m} c_m^* c_n \underbrace{\langle q_m | \hat{Q} | q_n \rangle}_{\text{still a sum}}$   
 $= \sum_{n,m} c_m^* c_n \langle q_m | \hat{Q} \sum_i |q_i\rangle \langle q_i | q_n \rangle$   
 $= \sum_{n,m,i} c_m^* c_n \sum_j \langle q_m | q_j \rangle \langle q_j | \hat{Q} | q_i \rangle \langle q_i | q_n \rangle$   
 $= \sum_{n,m,i,j} c_m^* c_n \langle q_m | q_j \rangle Q_{ij} \langle q_i | q_n \rangle$   
 $= \sum_{n,m,i,j} c_m^* c_n Q_{ij} \delta_{mj} \delta_{in}$

Q6: What if  $\hat{Q}$  has a continuous spectrum?  
 Use  $\mathbb{I} = \int |q\rangle\langle q| dq$  to show  $\langle \Psi | \hat{Q} | \Psi \rangle = \iint \langle \Psi | q' \rangle \langle q' | \hat{Q} | q \rangle \langle q | \Psi \rangle dq' dq$

$$\begin{aligned} \text{A6: } \langle \Psi | \hat{Q} | \Psi \rangle &= \langle \Psi | \hat{Q} \int |q\rangle\langle q| \Psi \rangle dq = \int \langle \Psi | \hat{Q} | q \rangle \langle q | \Psi \rangle dq \\ &= \int \left( \int \langle \Psi | q' \rangle \langle q' | \hat{Q} | q \rangle dq' \right) \langle q | \Psi \rangle dq \\ &= \iint \langle \Psi | q' \rangle \langle q' | \hat{Q} | q \rangle \langle q | \Psi \rangle dq dq' \end{aligned}$$

Q7: Show A6 reduces to  $\langle \Psi | \hat{x} | \Psi \rangle = \int \Psi^*(x) x \Psi(x) dx$   
 if  $\hat{Q} = \hat{x}$ .

$$\text{A7: } \langle \Psi | \hat{Q} | \Psi \rangle = \iint \langle \Psi | q' \rangle \langle q' | \hat{Q} | q \rangle \langle q | \Psi \rangle dq dq'$$

$$\text{Let } \hat{Q} = \hat{x}$$

$$\Rightarrow \langle \Psi | \hat{x} | \Psi \rangle = \iint \langle \Psi | x' \rangle \langle x' | \hat{x} | x \rangle \langle x | \Psi \rangle dx dx'$$

$$= \iint \langle \Psi | x' \rangle \langle x' | x | x \rangle \langle x | \Psi \rangle dx dx'$$

$$= \iint \langle \Psi | x' \rangle x \langle x' | x \rangle \langle x | \Psi \rangle dx dx'$$

$$= \iint \langle \Psi | x' \rangle x \delta(x' - x) \langle x | \Psi \rangle dx dx'$$

$$= \iint \Psi^*(x') x \delta(x' - x) \Psi(x) dx dx'$$

$$\langle \Psi | \hat{x} | \Psi \rangle = \int \Psi^*(x) x \Psi(x) dx$$

$$\begin{aligned} \langle \Psi | x' \rangle &= \Psi^*(x') \\ \langle x' | x \rangle &= \delta(x' - x) \\ \langle x | \Psi \rangle &= \Psi(x) \end{aligned}$$